

# Exact Volumes of Semi-Algebraic Convex Bodies

Lakshmi Ramesh<sup>b</sup> and Nicolas Weiss<sup>#</sup>

<sup>b</sup> University of Bielefeld. <sup>#</sup> MPI for Mathematics in the Sciences, Leipzig.



## Overview

We compute volumes of semi-algebraic convex bodies defined by finitely many concave polynomials with arbitrary precision. This is motivated by geometric statistics, where intersections of convex bodies arise as maximum likelihood estimator (MLE) sets [1].

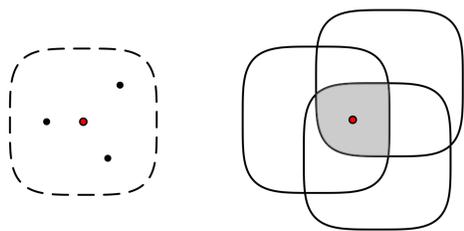


Figure: On the left: An  $\ell_4$ -ball with unknown center  $\theta$  in  $\mathbb{R}^2$  and  $p_1, \dots, p_k$  samples. On the right,  $K$  is translated to each of the  $p_i$ . The MLE set is the shaded region on the right.

Consider a convex body  $K + \theta$  with unknown center  $\theta$  and random samples  $X_1, \dots, X_k$  from  $K + \theta$ . The MLE set is then the bounded convex set

$$C = \bigcap_{i=1}^k (K + X_i).$$

One can think of the MLE set as an analog of the confidence interval in higher dimensions, i.e.

$$\text{vol}(C) = \text{uncertainty}(\mu \mid X_1, \dots, X_k).$$

We present an algorithm for the computation of the volume of these MLE sets based on the approach of [2]. The setting of concave polynomial, however, allows for various simplifications. We provide a full implementation of the algorithm in SageMath available on Github and Zenodo [3].

## Volumes as Periods of Rational Functions

Classically, the volume of semi-algebraic sets  $C$  is computed by *counting points* in  $C$ , for example with a Monte Carlo approach. This amounts to integrate the *volume form* over  $C$ . When  $C = \{f(x) > 0\}$  for  $f$  a smooth polynomial, one can also represent the volume as a period of a rational function:

$$\int_C 1 dx_1 \wedge \dots \wedge dx_n = \text{Vol}(C) = \frac{1}{2\pi i} \int_{\Gamma} \frac{x_1 \partial_1 \bullet f(x)}{f(x)} dx_1 \wedge \dots \wedge dx_n.$$

Whereas in the classical view  $C$  is an *open* and *real* integration contour in  $\mathbb{R}^n$ , the expression as a period of a rational function is an integral over a *closed* and *complex* contour  $\Gamma$  in  $\mathbb{C}^n$ . This contour is the *Leray coboundary* for  $\{x \in \mathbb{R}^n \mid f(x) = 0\}$ .

For  $f(t, x)$ , such that  $I(t) = \int_{\Gamma(t)} A(t, x) dx$ ,  $D$ -module theory naturally provides non-trivial differential equations for such parametric periods: Any element of the *integration ideal*

$$D_t \cap (\text{Ann}_{D_{t,x}}(A(t, x)) + \partial_{x_1} D_{t,x} + \dots + \partial_{x_n} D_{t,x}), \quad \text{where } D_{t,x} = \mathbb{Q}[t, \mathbf{x}] \langle \partial_t, \partial_{x_1}, \dots, \partial_{x_n} \rangle,$$

annihilates  $I(t)$ , provided that the contour  $\Gamma(t)$  is locally constant in a suitable sense. Integration ideals formalize *integration-by-parts* identities and are computed by a process called *creative telescoping*.

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## Volumes as Solutions of Differential Equations

Volumes of a convex semi-algebraic set  $C = \bigcap_i \{f_i > 0\}$  can be expressed as solutions to *univariate linear differential* equations in the following two ways. It relies on realizing  $C$  as one of the regions of  $\{\prod f_i > 0\}$ , therefore placing it in a family of smooth varieties defined by  $\prod f_i - t$ .

(A) - **Deformation:**

$$\text{Vol}(C) = \lim_{0 \leftarrow t} \text{Vol}(C_t)$$

(B) - **Integration:**

$$\text{Vol}(C_\epsilon) = \int_{a(\epsilon)}^{b(\epsilon)} \text{Vol}(C_\epsilon \cap \{x_1 = s\}) ds$$

The volume of the semi-algebraic set (or region thereof) is then realized as the limit respectively the integral of a holonomic function. And the linear differential operator  $P$  that annihilates it only depends on the defining polynomials  $f_1, \dots, f_k$ . If  $P$  has  $\text{order}(P) = r$ , then that holonomic function can be solved to *arbitrary precision* by *recursively* providing either  $r$  volumes of *deformations* respectively *lower-dimensional* slices.

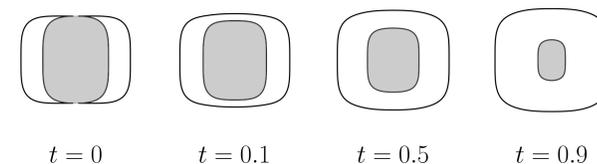


Figure: The intersection of two  $\ell_4$ -balls in  $\mathbb{R}^2$  defined by  $f_1 = 1 - x^4 - y^4$  and  $f_2 = 1 - (x - 1/2)^4 - y^4$  is deformed into the family of smoothly bounded sets  $C_t \subset \{\prod f_i - t > 0\}$ .  $\text{Vol}(C_t)$  is annihilated for  $t$  small by a linear differential operator  $P_t$  of order  $\text{ord}(P_t) = 6$  and degree  $\text{deg}(P_t) = 18$ .

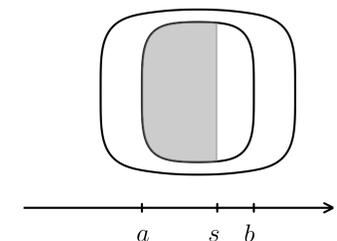


Figure: For  $f_1, f_2$  as on the left and for generic small  $t_0$ , the volume of the one-dimensional slices  $\text{Vol}(C_{t_0} \cap \{x = s\})$  is annihilated by a linear differential operator  $P_s$  of order  $\text{ord}(P_s) = 2$  and degree  $\text{deg}(P_s) = 23$ .

In the above example, the volume of the intersection is then determined by computing a total of 7 Picard–Fuchs operators and 12 one-dimensional slice volumes. Using high-precision solvers for linear ODEs one obtains to arbitrary precision  $\text{Vol}(C) = 2.708344826299720090001844936980643720528690610575117510163356615319 \dots$

## References

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